Beyond the Higgs Boson Discovery Energy Frontier Physics at the LHC

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University of Warwick Physics Colloquium 25 October 2017



Where we start: the Standard Model

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The Standard Model, SM





Force-carriers

Fermions

Bosons

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Higgs Discovery (ATLAS and CMS) July 4th 2012 (CERN and Melbourne)

10-8

110 115 120 125 130 135 140 145 150

Maximum excess observed at Local significance (including ener Probability of background up

Expected from SM Higgs mu=126.5

Global significance: 4.1-4.3 a (for LEE over 11)





~ 4



2 Ge/ ATLAS Data 3500 Excellent yy mass resolution crucial, as Sig+Bkg Fit (m_=126.5 GeV) Events / 3000 well as γ -ID to reject jet/ π^0 background Bkg (4th order polynomial) 2500 2000 1500E s=7 TeV, Ldt=4.8fb 1000E Is=8 TeV, ∫Ldt=5.9fb⁻¹ Η→γγ 500E (a) ate: 2011-10-22 15:30:29 UT 2 weights / 2 GeV Events - Bkg 200 100-100 -200 Data S/B Weighted 100 Sig+Bkg Fit (m_=126.5 GeV) Bkg (4th order polynomial) 80 60 40 Events weighted according to 20 (c) S/B in selected event category Σweights - Bkg Inclusive signal/background S/B ~3%

110

100

120

130

140

150

m_{yy} [GeV]

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160





 $H \rightarrow ZZ^* \rightarrow 4\ell$ "Golden channel" - excellent mass resolution and S/B~1





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Physics Letters B 716 (2012) 1-29

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Physics Letters B

www.elsevier.com/locate/physletb

Observation of a new particle in the search for the Standard Model Higgs boson with the ATLAS detector at the LHC \approx >7700 citations!

ATLAS Collaboration*

This paper is dedicated to the memory of our ATLAS colleagues who did not live to see the full impact and significance of their contributions to the experiment.

ARTICLE INFO

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ABSTRACT

A search for the Standard Model Higgs boson in proton–proton collisions with the ATLAS detector at the LHC is presented. The datasets used correspond to integrated luminosities of approximately 4.8 fb⁻¹ collected at $\sqrt{s} = 7$ TeV in 2011 and 5.8 fb⁻¹ at $\sqrt{s} = 8$ TeV in 2012. Individual searches in the channels $H \rightarrow ZZ^{(*)} \rightarrow 4\ell$, $H \rightarrow \gamma\gamma$ and $H \rightarrow WW^{(*)} \rightarrow ev\mu\nu$ in the 8 TeV data are combined with previously published results of searches for $H \rightarrow ZZ^{(*)}$, $WW^{(*)}$, $b\bar{b}$ and $\tau^+\tau^-$ in the 7 TeV data and results from improved analyses of the $H \rightarrow ZZ^{(*)} \rightarrow 4\ell$ and $H \rightarrow \gamma\gamma$ channels in the 7 TeV data. Clear evidence for the production of a neutral boson with a measured mass of 126.0 ± 0.4 (stat) ±0.4 (sys) GeV is presented. This observation, which has a significance of 5.9 standard deviations, corresponding to a background fluctuation probability of 1.7×10^{-9} , is compatible with the production and decay of the Standard Model Higgs boson.

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1. Introduction

120–135 GeV; using the existing LHC constraints, the observed local significances for $m_H = 125$ GeV are 2.7 σ for CDF [14], 1.1 σ for



The end of physics?

With the H discovery, we can identify all the particles in the SM with known states

But many questions are raised:

- Is it the SM Higgs boson?
- Is there only one scalar?
- Is the H(125) *solely* responsible for electroweak symmetry-breaking (EWSB)?
- Why is the H so light?

And we almost learned to stop asking the harder questions

- Why 3 generations of fermions?
- Why such *different* masses?
- Where is grand unification?
- What is dark matter & energy?
- Matter-antimatter asymmetry?
- Extra dimensions? Branes?

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Forces

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Precision measurements of the Higgs sector, and of EWSB in general, are mandated

Quarks

Searches for new physics are required to try to resolve these - but we do not know the energy scale of such beyond-SM physics, except for general indications that it should be in the TeV-scale range

The LHC and experiments

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Proton-proton & heavy-ion collisions

Lake Geneva

CMS

RANCE

LHC ring: 27 km circumference ~100 m underground HCh

LHC 27 km

Airpor

CERN Prévessin

10 V

ATLAS

CERN main site

ALICE







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ATLAS

HCb Airport

CERN Prévessin

-

LHC 27 km

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Proton-proton & heavy-ion collisions

Lake Geneva

CMS



LHCb Airport

-

CERN Prévessin

LHC 27 km

CERN main site

ALICE

LHC ring: 27 km circumference ~100 m underground





Detector principles





Multiple detector layers measure charged particle momenta (tracks), EM and hadronic energies (calorimetry), and provide particle identification from different signatures *Full event: missing transverse momentum balance can be used* \rightarrow *sensitive to invisible particles (v, new physics - dark matter?)*

ATLAS Detector

7000t, 45m long x 25m diameter Si+transition radiation tracker, 2T solenoid, LAr sampling calorimetry, large air-core muon system



~110 M channels, with timing capable of separating particles from adjacent bunch-crossings (25ns) Physical size allows precise momentum measurements and provides material depth to absorb TeV-energy jets of hadrons



Adelaide, Albany, Alberta, NIKHEF Amsterdam, Ankara, LAPP Annecy, Argonne NL, Arizona, UT Arlington, Athens, NTU Athens, UT Austin, Baku, IFAE Barcelona, Belgrade, Bergen, Berkeley LBL and UC, HU Berlin, Bern, Birmingham, UAN Bogota, Bologna, Bonn, Boston, Brandeis, Bratislava/SAS Kosice, Brazil Cluster, Brookhaven NL, Buenos Aires, Bucharest, Cambridge, Carleton, CERN, China IHEP-NJU-THU, China USTC-SDU-SJTU, Chicago, Chile, Clermont-Ferrand, Columbia, NBI Copenhagen, Cosenza, AGH UST Cracow, IFJ PAN Cracow, SMU Dallas, UT Dallas, Demokritos, DESY, Dortmund, TU Dresden, JINR Dubna, Duke, Edinburgh, Frascati, Freiburg, Geneva, Genoa, Giessen, Glasgow, Göttingen, LPSC Grenoble, Technion Haifa, Harvard, Heidelberg, Hiroshima IT, Hong Kong, NTHU Hsinchu, Indiana, Innsbruck, Iowa SU, Iowa, UC Irvine, Istanbul Bogazici, KEK, Kobe, Kyoto, Kyoto UE, Kyushu, Lancaster, UN La Plata, Lecce, Lisbon LIP, Liverpool, Ljubljana, QM London, RH London, UC London, Louisiana Tech, Lund, UA Madrid, Mainz, Manchester, CPPM Marseille, Massachusetts, MIT, Melbourne, Michigan, Michigan SU, Milano, Minsk NAS, Minsk NCPHEP, Montreal, McGill Montreal, RUPHE Morocco, FIAN Moscow, ITEP Moscow, MEPH Moscow, MSU Moscow, Munich LMU, MPI Munich, Nagasaki IAS, Nagoya, Naples, New Mexico, New York, Nijmegen, Northern Illinois University, Novosibirsk BINP-NSU, NPI Petersburg, Ohio SU, Okayama, Oklahoma, Oklahoma, SU, Olomouc, Oregon, LAL Orsay, Osaka, Oslo, Oxford, Paris VI and VII, Pavia, Pennsylvania, Pisa, Pittsburgh, CAS Prague, CU Prague, TU Prague, IHEP Protvino, Rome I, Rome II, Rome III, RAL-STFC, DAPNIA Saclay, Santa Cruz UC, Sheffield, Shinshu, Siegen, Simon Fraser Burnaby, SLAC, South Africa Cluster, Stockholm, KTH Stockholm, Stony Brook, Sydney, Sussex, AS Taipei, Tbilisi, Tel Aviv, Thessaloniki, Tokyo ICEPP, Tokyo MU, Tokyo Tech, Tomsk SU, Toronto, Trento, TRIUMF, Tsukuba, Tufts, Udine/ICTP, Uppsala, UI Urbana, Valencia, UBC Vancouver, Victoria, Warwick, Waseda, Washington, Weizmann Rehovot, FH Wiener Neustadt, Wisconsin, Wu

The first decades of LHC operation



Why we went to 13 TeV

"Interesting" processes come from hard scattering of partons

- Colliding *partons* carry a fraction of the proton momentum, x_i , according to a parton density function ("pdf")
- The partonic centre-of-mass energy is

$$\sqrt{\hat{s}} = \sqrt{x_1 x_2 s} = M_X$$

 $\sigma =$



 M_{x} (GeV)



Ratio of parton-parton luminosity for pp centre-

 $\left(\frac{d L(a,b)}{d \hat{s}}\right)$

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Why we went to 13 TeV







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The first decades of LHC operation



The first decades of LHC operation



2015 - the restart at 13 TeV



Integrated luminosity $\int Ldt$ drives the signal event yield N_{obs}

 $= \sigma \epsilon_{exp} \int Ldt$ obs 25 Integrated Luminosity [fb⁻¹] 2012 σ : cross-section $\sigma(\sqrt{s})$ ε_{exp} : experimental efficiency 20 15 Almost like starting a brand new accelerator: a late 10 start in the year, and 2011 various problems after the two-year stop A slow year... 2015

02-May

01-Jul

31-Aug

02-Mar

31-Dec

31-Oct

2016 - a great production year



Integrated luminosity $\int Ldt$ drives the signal event yield N_{obs}



2017 - another good year



Integrated luminosity $\int Ldt$ drives the signal event yield N_{obs}



Month in Year



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The trigger challenge

Complex trigger menu designed to meet varied physics, monitoring and performance requirements

- Typically ~2000 active menu items
- Key is to keep stable, well understood, main primary triggers
- Use as low energy thresholds as possible to keep efficiencies as high as possible
- Small and active community manages trigger menus to optimise physics output



- $E_{\tau}(e) > 24-26$
- *p*₇(µ) > 24-26
- $E_{T}^{\text{miss}} > 90-110$
- $E_{\tau}(jet) > 380$
- $E_{\tau}(\gamma) > 140$
- $p_{\tau}(\mu 1, \mu 2) > 6,6 + \text{topo/mass selections}$
- $E_{\tau}(\gamma 1, \gamma 2) > 35, 25$









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H(125) production & decay



A 125 GeV Higgs boson is a convenient object experimentally - many production and decay modes should be measurable

- Production and decay processes probe couplings of H to different particles
- Is it the Standard Model Higgs or not?



H(125) production modes

Combined analysis of Run-1 data: H(125) production & decays

With assumptions about decays, we can probe the different production processes (normalised rates " μ " (=1 in SM))



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- Able to separate statistically the ggF and VBF processes
- Not yet VH or ttH at 5σ
- Observing ttH production is a key Run-2 goal

These are not yet precision measurements but few percent errors should be obtainable with the expected LHC samples

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H(125) production & decay



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H(125) decay modes

Combining ATLAS and CMS Run-1 data, observed (at $>5\sigma$ significance)

- Н→үү
- H→ZZ(→4ℓ (ℓ=e,µ))
- H→WW
- Η→ττ

Run-1 data not yet sensitive to the dominant $H\rightarrow$ bb, or rarer, decays, such as to second generation fermions μ , c, s

No direct evidence of Higgs coupling to (any) quarks from Run-1 data Coupling to only one lepton species (τ) observed



Decay signal strengths relative to Standard Model "µ" (=1 in SM)





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Mass of the H("125")



Recall that m_{μ} is a free parameter in the Standard Model

- To measure $m_{_{\!H}}$, we use $\gamma\gamma$ and 4 ℓ decays, where we can reconstruct the mass event-by-event with high resolution
- Requires excellent understanding of energy scales for lepton/photons



Calibrate detector performance relative to simulations using very large and clean samples of decays of particles of known mass, here: $J/\psi, \Upsilon, Z \rightarrow ee/\mu\mu$

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Already a precision measurement: 2 per-mille relative error dominated by statistical not systematic uncertainties

Measuring H(125) at 13 TeV in Run-2





Clear signals in $\gamma\gamma$ and 4 ℓ \rightarrow combined $\sigma(pp \rightarrow H)$ at 13 TeV \rightarrow overall significance at 13 TeV ~10 σ



First evidence for H→bb

Hunt for $H \rightarrow b\overline{b}$ decay, produced with a W or a Z boson

• H \rightarrow bb dominant Higgs decay in SM: BR~58%





Signal strength relative to SM expected value:

 $0.90 \pm 0.18(\text{stat.})^{+0.21}_{-0.19}(\text{syst.})$

Significance 3.5σ (expected 3.0σ)

First evidence for H decay to quarks - observing this at 5 remains a key goal for Run-2

$H \rightarrow \mu \mu$ - rare decay: 2nd generation!





With full 2015+2016 data, look for a peak in dimuon mass spectrum

• Event categories improve sensitivity

No excess observed \rightarrow place limits on signal strength μ_s relative to Standard Model, combine also with (weaker) Run-1 results:

 μ_{s} < 2.8 at 95% CL (2.9 expected)



$H \rightarrow \mu \mu$ - rare decay: 2nd generation!





Still a way to go here - precise measurements of Higgs couplings to any second generation fermion will require much more data (\rightarrow HL-LHC)

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With full 2015+2016 data, look for a peak in dimuon mass spectrum

• Event categories improve sensitivity

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 μ_{s} < 2.8 at 95% CL (2.9 expected)



48





W-boson mass





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Electroweak precision test





measurements of W and top masses

Within the SM framework, the precise value of m_w is related to other quantities via:

$$m_W^2 \left(1 - \frac{m_W^2}{m_Z^2} \right) = \frac{\pi \alpha}{\sqrt{2}G_\mu} (1 + \Delta r),$$

Where Δr includes radiative effects (loops), and so depends on (e.g.) $m_{_H}$ and $m_{_{top}}$



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Massive diboson production

Run-1 puzzle to describe inclusive diboson cross-sections

Measurements tended to lie above NLO calculations

NNLO calculations \rightarrow ~20% corrections and better agreement



 $\overline{q'}$

Standard Model Production Cross Section Measurements

Status: July 2017



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a.k.a. Beyond-the-SM, BSM, searches

There are far too many to describe (hundreds of results) - I pick just one of my favourites

Dijet resonance search

Search for new physics in dijet invariant mass spectrum



Examples (at 95% CL):

m(q*) > 6.0 TeV (Cf Run-1: 4.1 TeV)

m(W*) > 3.4 TeV and not within 3.77-3.85 TeV



arxiv:1703.09127



Highest-mass central dijet event (2016) - m(jj)=8.2 TeV



Run: 305777 Event: 4144227629 2016-08-08 08:51:15 CEST

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ATLAS Exotics Searches* - 95% CL Upper Exclusion Limits

Status: July 2017

Status: July 2017 $\int \mathcal{L} dt = (3.2 - 37.0) \text{ fb}^{-1} \sqrt{s}$									
	Model	<i>ℓ</i> ,γ	Jets†	$\mathbf{E}_{\mathrm{T}}^{\mathrm{miss}}$	∫£ dt[fb	-1] Limit		Reference	
Extra dimensions	ADD $G_{KK} + g/q$ ADD non-resonant $\gamma\gamma$ ADD QBH ADD BH high $\sum p_T$ ADD BH multijet RS1 $G_{KK} \rightarrow \gamma\gamma$ Bulk RS $G_{KK} \rightarrow WW \rightarrow qq\ell\nu$ 2UED / RPP	$\begin{array}{c} 0 \ e, \mu \\ 2 \ \gamma \\ - \\ \geq 1 \ e, \mu \\ - \\ 2 \ \gamma \\ 1 \ e, \mu \\ 1 \ e, \mu \end{array}$	1 - 4j - 2 j $\ge 2j$ $\ge 3j$ - 1 J $\ge 2b, \ge 3$	Yes – – – – Yes j Yes	36.1 36.7 37.0 3.2 3.6 36.7 36.1 13.2	Mp 7.75 TeV Ms 8.6 TeV Mth 8.9 TeV Mth 8.2 TeV Mth 9.55 TeV GKK mass 4.1 TeV KK mass 1.75 TeV KK mass 1.6 TeV	$\begin{split} n &= 2\\ n &= 3 \text{ HLZ NLO}\\ n &= 6\\ n &= 6, M_D = 3 \text{ TeV, rot BH}\\ n &= 6, M_D = 3 \text{ TeV, rot BH}\\ k/\overline{M}_{Pl} &= 0.1\\ k/\overline{M}_{Pl} &= 1.0\\ \text{Tier } (1,1), \mathcal{B}(\mathcal{A}^{(1,1)} \rightarrow tt) = 1 \end{split}$	ATLAS-CONF-2017-060 CERN-EP-2017-132 1703.09217 1606.02265 1512.02586 CERN-EP-2017-132 ATLAS-CONF-2017-051 ATLAS-CONF-2016-104	
Gauge bosons	$\begin{array}{l} \mathrm{SSM} \ Z' \to \ell\ell \\ \mathrm{SSM} \ Z' \to \tau\tau \\ \mathrm{Leptophobic} \ Z' \to bb \\ \mathrm{Leptophobic} \ Z' \to tt \\ \mathrm{SSM} \ W' \to \ell\nu \\ \mathrm{HVT} \ V' \to WV \to qqqq \ \mathrm{model} \ \mathrm{B} \\ \mathrm{HVT} \ V' \to WH/ZH \ \mathrm{model} \ \mathrm{B} \\ \mathrm{LRSM} \ W'_R \to tb \\ \mathrm{LRSM} \ W'_R \to tb \end{array}$	$2 e, \mu$ 2τ $-$ $1 e, \mu$ $1 e, \mu$ $0 e, \mu$ ulti-channa $1 e, \mu$ $0 e, \mu$	- 2 b ≥ 1 b, ≥ 1J/ - 2 J el 2 b, 0-1 j ≥ 1 b, 1 J	- - 2j Yes Yes - Yes -	36.1 36.1 3.2 3.2 36.1 36.7 36.1 20.3 20.3	Z' mass 4.5 TeV Z' mass 2.4 TeV Z' mass 1.5 TeV Z' mass 2.0 TeV W' mass 5.1 TeV V' mass 3.5 TeV V' mass 2.93 TeV W' mass 1.92 TeV W' mass 1.76 TeV	$\Gamma/m = 3\%$ $g_V = 3$ $g_V = 3$	ATLAS-CONF-2017-027 ATLAS-CONF-2017-050 1603.08791 ATLAS-CONF-2016-014 1706.04786 CERN-EP-2017-147 ATLAS-CONF-2017-055 1410.4103 1408.0886	
CI	Cl qqqq Cl ℓℓqq Cl uutt 20	– 2 e,μ (SS)/≥3 e,	2 j _ µ ≥1 b, ≥1 j	– – Yes	37.0 36.1 20.3	Λ Λ Λ 4.9 TeV	21.8 TeV η_{LL}^- 40.1 TeV η_{LL}^- $ C_{RR} = 1$	1703.09217 ATLAS-CONF-2017-027 1504.04605	
MQ	Axial-vector mediator (Dirac DM) Vector mediator (Dirac DM) $VV_{\chi\chi}$ EFT (Dirac DM)	0 e, μ 0 e, μ, 1 γ 0 e, μ	1 - 4 j $\leq 1 j$ $1 J, \leq 1 j$	Yes Yes Yes	36.1 36.1 3.2	m _{med} 1.5 TeV m _{med} 1.2 TeV M _* 700 GeV	$\begin{array}{l} g_q{=}0.25,g_\chi{=}1.0,m(\chi)<400~{\rm GeV}\\ g_q{=}0.25,g_\chi{=}1.0,m(\chi)<480~{\rm GeV}\\ m(\chi)<150~{\rm GeV} \end{array}$	ATLAS-CONF-2017-060 1704.03848 1608.02372	
ГQ	Scalar LQ 1 st gen Scalar LQ 2 nd gen Scalar LQ 3 rd gen	2 e 2 μ 1 e, μ	≥ 2 j ≥ 2 j ≥1 b, ≥3 j	– – Yes	3.2 3.2 20.3	LQ mass 1.1 TeV LQ mass 1.05 TeV LQ mass 640 GeV	$egin{array}{ll} eta = 1 \ eta = 1 \ eta = 1 \ eta = 0 \end{array}$	1605.06035 1605.06035 1508.04735	
Heavy quarks	$ \begin{array}{l} VLQ \ TT \rightarrow Ht + X \\ VLQ \ TT \rightarrow Zt + X \\ VLQ \ TT \rightarrow Wb + X \\ VLQ \ BB \rightarrow Hb + X \\ VLQ \ BB \rightarrow Zb + X \\ VLQ \ BB \rightarrow Wt + X \\ VLQ \ QQ \rightarrow WqWq \end{array} $	0 or 1 <i>e</i> , <i>µ</i> 1 <i>e</i> , <i>µ</i> 1 <i>e</i> , <i>µ</i> 2/≥3 <i>e</i> , <i>µ</i> 1 <i>e</i> , <i>µ</i> 1 <i>e</i> , <i>µ</i>	$\begin{array}{l} \geq 2 \ b, \geq 3 \\ \geq 1 \ b, \geq 3 \\ \geq 1 \ b, \geq 1 J / \\ \geq 2 \ b, \geq 3 \\ \geq 2 / \geq 1 \ b \\ \geq 1 \ b, \geq 1 J / \\ \geq 4 \ j \end{array}$	j Yes j Yes 2j Yes j Yes – 2j Yes Yes	13.2 36.1 20.3 20.3 36.1 20.3	T mass1.2 TeVT mass1.16 TeVT mass1.35 TeVB mass700 GeVB mass790 GeVB mass1.25 TeVQ mass690 GeV	$\begin{split} \mathcal{B}(T \to Ht) &= 1\\ \mathcal{B}(T \to Zt) &= 1\\ \mathcal{B}(T \to Wb) &= 1\\ \mathcal{B}(B \to Hb) &= 1\\ \mathcal{B}(B \to Zb) &= 1\\ \mathcal{B}(B \to Wt) &= 1 \end{split}$	ATLAS-CONF-2016-104 1705.10751 CERN-EP-2017-094 1505.04306 1409.5500 CERN-EP-2017-094 1509.04261	
Excited fermions	Excited quark $q^* \rightarrow qg$ Excited quark $q^* \rightarrow q\gamma$ Excited quark $b^* \rightarrow bg$ Excited quark $b^* \rightarrow Wt$ Excited lepton ℓ^* Excited lepton γ^*	- 1 γ - 1 or 2 e, μ 3 e, μ 3 e, μ, τ	2 j 1 j 1 b, 1 j 1 b, 2-0 j - -	- - Yes -	37.0 36.7 13.3 20.3 20.3 20.3	q* mass 6.0 TeV q* mass 5.3 TeV b* mass 2.3 TeV b* mass 1.5 TeV /* mass 3.0 TeV /* mass 1.6 TeV	only u^* and d^* , $\Lambda = m(q^*)$ only u^* and d^* , $\Lambda = m(q^*)$ $f_g = f_L = f_R = 1$ $\Lambda = 3.0 \text{ TeV}$ $\Lambda = 1.6 \text{ TeV}$	1703.09127 CERN-EP-2017-148 ATLAS-CONF-2016-060 1510.02664 1411.2921 1411.2921	
Other	LRSM Majorana ν Higgs triplet $H^{\pm\pm} \rightarrow \ell \ell$ 2,5 Higgs triplet $H^{\pm\pm} \rightarrow \ell \tau$ Monotop (non-res prod) Multi-charged particles Magnetic monopoles	2 e, μ 3,4 e, μ (S: 3 e, μ, τ 1 e, μ - - -	2 j S) - - 1 b - -	- - Yes - - 3 TeV	20.3 36.1 20.3 20.3 20.3 7.0	N° mass 2.0 TeV H** mass 870 GeV H** mass 400 GeV spin-1 invisible particle mass 657 GeV multi-charged particle mass 785 GeV monopole mass 1.34 TeV 10 ⁻¹ 1	$m(W_R) = 2.4 \text{ TeV, no mixing}$ DY production DY production, $\mathcal{B}(H_L^{\pm\pm} \rightarrow \ell \tau) = 1$ $a_{\text{non-res}} = 0.2$ DY production, $ q = 5e$ DY production, $ g = 1g_D$, spin 1/2 Mass scale [TeV]	1506.06020 ATLAS-CONF-2017-053 1411.2921 1410.5404 1504.04188 1509.08059	

ATLAS Preliminary

*Only a selection of the available mass limits on new states or phenomena is shown.

†Small-radius (large-radius) jets are denoted by the letter j (J).

ATLAS Exotics Searches* - 95% CL Upper Exclusion Limits

Status: Julv 2017



ATLAS Preliminary

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SUSY Searches

EXPERIM

ATLAS Preliminary

ATLAS SUSY Searches* - 95% CL Lower Limits May 2017

May 2017 $\sqrt{s} = 7, 8, 13$ TeV									
	Model	e, μ, τ, γ	Jets	E ^{miss} T	∫ <i>L dt</i> [fb	⁻¹] Mass limit	$\sqrt{s} = 7, 8$	TeV $\sqrt{s} = 13$ TeV	Reference
Inclusive Searches	$ \begin{array}{l} \text{MSUGRA/CMSSM} \\ \tilde{q}\tilde{q}, \; \tilde{q} \rightarrow q \tilde{\chi}_{1}^{0} \\ \tilde{q}\tilde{q}, \; \tilde{q} \rightarrow q \tilde{\chi}_{1}^{1} \; (\text{compressed}) \\ \tilde{g}\tilde{s}, \; \tilde{g} \rightarrow q \tilde{\chi}_{1}^{1} \; (\text{compressed}) \\ \tilde{g}\tilde{s}, \; \tilde{g} \rightarrow q q \tilde{\chi}_{1}^{1} \\ \tilde{g}\tilde{s}, \; \tilde{g} \rightarrow q q \mathcal{K}_{1}^{1} \rightarrow q q \Psi^{\pm} \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{s}, \; \tilde{g} \rightarrow q q (\mathcal{E}/\gamma) \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{s}, \; \tilde{g} \rightarrow q q \mathcal{W} Z \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{s}, \; \tilde{g} \rightarrow q q \mathcal{W} Z \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{s}, \; \tilde{g} \rightarrow q q \mathcal{W} Z \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{s}, \; \tilde{g} \rightarrow q q \mathcal{W} Z \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{s}, \; \tilde{g} \rightarrow q q \mathcal{W} Z \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{s}, \; \tilde{g} \rightarrow q q \mathcal{W} Z \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{s}, \; \tilde{s} \rightarrow q q \mathcal{W} Z \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{s}, \; \tilde{s} \rightarrow q q \mathcal{W} Z \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{s}, \; \tilde{s} \rightarrow q q \mathcal{W} Z \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{s}, \; \tilde{s} \rightarrow q q \mathcal{W} Z \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{s}, \; \tilde{s} \rightarrow q q \mathcal{W} Z \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{s}, \; \tilde{s} \rightarrow q q \mathcal{W} Z \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{s}, \; \tilde{s} \rightarrow q q \mathcal{W} Z \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{s}, \; \tilde{s} \rightarrow q q \mathcal{W} Z \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{s}, \; \tilde{s} \rightarrow q q \mathcal{W} Z \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{s}, \; \tilde{s} \rightarrow q q \mathcal{W} Z \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{s}, \; \tilde{s} \rightarrow q q \mathcal{W} Z \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{s}, \; \tilde{s} \rightarrow q q \mathcal{W} Z \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{s}, \; \tilde{s} \rightarrow q q \mathcal{W} Z \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{s}, \; \tilde{s} \rightarrow q q \mathcal{W} Z \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{s}, \; \tilde{s} \rightarrow q q \mathcal{W} Z \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{s}, \; \tilde{s} \rightarrow q q \mathcal{W} Z \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{s}, \; \tilde{s} \rightarrow q q \mathcal{W} Z \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{s}, \; \tilde{s} \rightarrow q \mathcal{W} Z \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{s}, \; \tilde{s} \rightarrow q \mathcal{W} Z \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{s}, \; \tilde{s} \rightarrow q \mathcal{W} Z \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{s}, \; \tilde{s} \rightarrow q \mathcal{W} Z \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{s}, \; \tilde{s} \rightarrow q \mathcal{W} Z \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{s}\tilde{s}, \; \tilde{s} \rightarrow q \mathcal{W} Z \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{s}\tilde{s}, \; \tilde{s} \rightarrow q \mathcal{W} Z \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{s}\tilde{s}, \; \tilde{s} \rightarrow q \mathcal{W} Z \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{s}\tilde{s}\tilde{s}\tilde{s}\tilde{s} \tilde{s} \tilde{s} \tilde{s} s$	$\begin{array}{c} 0\text{-}3 \ e, \mu/1\text{-}2 \ \tau \\ 0 \\ \text{mono-jet} \\ 0 \\ 3 \ e, \mu \\ 0 \\ 1\text{-}2 \ \tau + 0\text{-}1 \ \ell \\ 2 \ \gamma \\ \gamma \\ 2 \ e, \mu \ (Z) \\ 0 \end{array}$	2-10 jets/3 <i>b</i> 2-6 jets 1-3 jets 2-6 jets 2-6 jets 4 jets 7-11 jets 0-2 jets 2 jets 2 jets 2 jets 2 jets	Yes Yes Yes Yes Yes Yes Yes Yes Yes Yes	20.3 36.1 3.2 36.1 36.1 36.1 3.2 3.2 20.3 13.3 20.3 20.3	\$\vec{a}{2}\$	1.85 TeV 1.57 TeV 2.02 TeV 2.01 TeV 1.825 TeV 1.8 TeV 2.0 TeV 1.65 TeV 3.7 TeV 1.8 TeV	$\begin{split} & m(\tilde{q}) \!=\! m(\tilde{g}) \\ & m(\tilde{x}_1^0) \! <\! 200 \; GeV, \; m(1^{st} \; gen, \tilde{q}) \!=\! m(2^{nd} \; gen, \tilde{q}) \\ & m(\tilde{x}_1^0) \! <\! 200 \; GeV \\ & m(\tilde{x}_1^0) \! <\! 400 \; GeV \\ & cr(NLSP) \! <\! 0.1 \; mm \\ & m(\tilde{x}_1^0) \! <\! 950 \; GeV, \; cr(NLSP) \! <\! 0.1 \; mm, \; \mu \! <\! 0 \\ & m(\tilde{x}_1^0) \! <\! 950 \; GeV, \; cr(NLSP) \! <\! 0.1 \; mm, \; \mu \! >\! 0 \\ & m(NLSP) \! >\! \! >\! 30 \; GeV \\ & m(\tilde{c}) \! >\! 1.8 \times 10^{-4} \; eV, \; m(\tilde{g}) \! =\!\! n(\tilde{q}) \! =\!\! 1.5 \; TeV \end{split}$	1507.05525 ATLAS-CONF-2017-022 1604.07773 ATLAS-CONF-2017-022 ATLAS-CONF-2017-022 ATLAS-CONF-2017-030 ATLAS-CONF-2017-033 1607.05979 1606.09150 1507.05493 ATLAS-CONF-2016-066 1503.03290 1502.01518
3 rd gen. ẽ med.	$\begin{array}{c} \tilde{g}\tilde{g}, \tilde{g} \rightarrow b \bar{b} \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{g}, \tilde{g} \rightarrow t \tilde{\mathcal{X}}_{1}^{0} \\ \tilde{g}\tilde{g}, \tilde{g} \rightarrow b \tilde{\mathcal{X}}_{1}^{+} \end{array}$	0 0-1 <i>e</i> ,μ 0-1 <i>e</i> ,μ	3 b 3 b 3 b	Yes Yes Yes	36.1 36.1 20.1	ž ž ž 1.3	1.92 TeV 1.97 TeV .37 TeV	$m(\tilde{\chi}_{1}^{0}) < 600 \text{ GeV} m(\tilde{\chi}_{1}^{0}) < 200 \text{ GeV} m(\tilde{\chi}_{1}^{0}) < 300 \text{ GeV}$	ATLAS-CONF-2017-021 ATLAS-CONF-2017-021 1407.0600
3 rd gen. squarks direct production	$ \begin{split} \bar{b}_1 \bar{b}_1, \bar{b}_1 \to b \bar{k}_1^0 \\ \bar{b}_1 \bar{b}_1, \bar{b}_1 \to b \bar{\ell}_1^{\bar{\pi}} \\ \bar{t}_1 \bar{t}_1, \bar{t}_1 \to b \bar{\ell}_1^{\bar{\pi}} \\ \bar{t}_1 \bar{t}_1, \bar{t}_1 \to b \bar{\ell}_1^{\bar{\pi}} \\ \bar{t}_1 \bar{t}_1, \bar{t}_1 \to b \bar{\ell}_1^0 \\ \bar{t}_1 \bar{t}_1, \bar{t}_1 \to c \bar{\ell}_1^0 \\ \bar{t}_1 \bar{t}_1 (natural GMSB) \\ \bar{t}_2 \bar{t}_2, \bar{t}_2 \to \bar{t}_1 + Z \\ \bar{t}_2 \bar{t}_2, \bar{t}_2 \to \bar{t}_1 + h \end{split} $	$\begin{matrix} 0 \\ 2 e, \mu (SS) \\ 0-2 e, \mu \\ 0-2 e, \mu \\ 0 \\ 2 e, \mu (Z) \\ 3 e, \mu (Z) \\ 1-2 e, \mu \end{matrix}$	2 b 1 b 1-2 b 0-2 jets/1-2 b mono-jet 1 b 1 b 4 b	Yes Yes Yes 4 Yes 20 Yes Yes Yes Yes	36.1 36.1 7/13.3).3/36.1 3.2 20.3 36.1 36.1	\$\vec{b}_1\$ 950 GeV \$\vec{b}_1\$ 275-700 GeV \$\vec{l}_1\$ 117-170 GeV 200-720 GeV \$\vec{l}_1\$ 90-198 GeV 205-950 GeV \$\vec{l}_1\$ 90-323 GeV 150-600 GeV \$\vec{l}_2\$ 290-790 GeV \$\vec{l}_2\$ \$\vec{l}_2\$ 320-880 GeV \$\vec{l}_2\$		$\begin{split} & m(\tilde{k}_1^0)\!<\!\!420~GeV \\ & m(\tilde{k}_1^0)\!<\!\!200~GeV, m(\tilde{k}_1^+)\!=m(\tilde{k}_1^0)\!+\!100~GeV \\ & m(\tilde{k}_1^+)\!=\!2m(\tilde{k}_1^0), m(\tilde{k}_1^0)\!=\!55~GeV \\ & m(\tilde{k}_1^0)\!=\!1~GeV \\ & m(\tilde{k}_1^0)\!=\!5~GeV \\ & m(\tilde{k}_1^0)\!>\!150~GeV \\ & m(\tilde{k}_1^0)\!=\!0~GeV \\ & m(\tilde{k}_1^0)\!=\!0~GeV \end{split}$	ATLAS-CONF-2017-038 ATLAS-CONF-2017-030 1209.2102, ATLAS-CONF-2016-077 1506.08616, ATLAS-CONF-2017-020 1604.07773 1403.5222 ATLAS-CONF-2017-019 ATLAS-CONF-2017-019
EW direct	$ \begin{split} \tilde{\ell}_{L,R} \tilde{\ell}_{L,R}, \tilde{\ell} \to \ell \tilde{\chi}_1^0 \\ \tilde{\chi}_1^+ \tilde{\chi}_1^-, \tilde{\chi}_1^+ \to \tilde{\ell}_V(\ell \tilde{\nu}) \\ \tilde{\chi}_1^+ \tilde{\chi}_1^-, \tilde{\chi}_1^+ \to \tilde{\ell}_V(\ell \tilde{\nu}), \\ \tilde{\chi}_1^+ \tilde{\chi}_2^0 \to \tilde{\ell}_L, \tilde{\nu}_L(\ell \tilde{\nu}), \ell \tilde{\nu}_L \ell(\tilde{\nu} \nu) \\ \tilde{\chi}_1^+ \tilde{\chi}_2^0 \to W_V^{01} 2 \chi_1^0 \\ \tilde{\chi}_1^+ \tilde{\chi}_2^0 \to W_V^{01} h \tilde{\chi}_1^+, h \to b \tilde{b} / W W / \tau \tau / \tau \\ \tilde{\chi}_2^0 \tilde{\chi}_3, \tilde{\chi}_{2,3}^0 \to \tilde{\ell}_R \ell \\ GGM (bino NLSP) weak prod., \tilde{\chi} \\ GGM (bino NLSP) weak prod., \tilde{\chi} \end{split} $	$2 e, \mu$ $2 e, \mu$ 2τ $3 e, \mu$ $2 \cdot 3 e, \mu$ $\gamma \gamma e, \mu, \gamma$ $4 e, \mu$ $\gamma^{0} \rightarrow \gamma \widetilde{G} 1 e, \mu + \gamma$ $0 \rightarrow \gamma \widetilde{G} 2 \gamma$	0 0 0-2 jets 0-2 b 0 -	Yes Yes Yes Yes Yes Yes Yes Yes	36.1 36.1 36.1 36.1 20.3 20.3 20.3 20.3	$ \begin{array}{c c c c c c c c c c c c c c c c c c c $	eV m (\tilde{x}_1^*) =m (\tilde{x}_2^0) =m	$\begin{split} & m(\tilde{\xi}_{1}^{0}) {=} 0 \\ & m(\tilde{\xi}_{1}^{0}) {=} 0, m(\tilde{\ell}, \tilde{\gamma}) {=} 0. 5(m(\tilde{\xi}_{1}^{+}) {+} m(\tilde{\xi}_{1}^{0})) \\ & m(\tilde{\xi}_{1}^{0}) {=} 0, m(\tilde{\ell}, \tilde{\gamma}) {=} 0. 5(m(\tilde{\xi}_{1}^{+}) {+} m(\tilde{\xi}_{1}^{0})) \\ & (\tilde{\xi}_{2}^{0}), m(\tilde{\xi}_{1}^{0}) {=} 0, m(\tilde{\ell}, \tilde{\gamma}) {=} 0. 5(m(\tilde{\xi}_{1}^{0}) {+} m(\tilde{\xi}_{1}^{0})) \\ & m(\tilde{\xi}_{1}^{0}) {=} m(\tilde{\xi}_{2}^{0}), m(\tilde{\xi}_{1}^{0}) {=} 0, \tilde{\ell} \text{ decoupled} \\ & m(\tilde{\xi}_{1}^{0}) {=} m(\tilde{\xi}_{2}^{0}), m(\tilde{\ell}, \tilde{\gamma}) {=} 0. 5(m(\tilde{\xi}_{2}^{0}) {+} m(\tilde{\xi}_{1}^{0})) \\ & cr {<} 1 mm \\ & cr {<} 1 mm \end{split}$	ATLAS-CONF-2017-039 ATLAS-CONF-2017-039 ATLAS-CONF-2017-035 ATLAS-CONF-2017-039 ATLAS-CONF-2017-039 1501.07110 1405.5086 1507.05493
Long-lived particles	$\begin{array}{l} \text{Direct} \tilde{X}_1^{\dagger} \tilde{X}_1^{-} \text{ prod., long-lived } \tilde{X}_1^{\pm} \\ \text{Direct} \tilde{X}_1^{\dagger} \tilde{X}_1^{-} \text{ prod., long-lived } \tilde{X}_1^{\pm} \\ \text{Stable, stopped } \tilde{g} \text{ R-hadron} \\ \text{Stable } \tilde{g} \text{ R-hadron} \\ \text{Metastable } \tilde{g} \text{ R-hadron} \\ \text{GMSB, stable } \tilde{\tau}, \tilde{X}_1^{0} \rightarrow \tilde{\tau}(\tilde{e}, \tilde{\mu}) + \tau(e, \\ \text{GMSB}, \tilde{X}_1^{0} \rightarrow \gamma \tilde{G}, \text{ long-lived } \tilde{X}_1^{0} \\ \tilde{g} \tilde{g}, \tilde{X}_1^{0} \rightarrow e e e^j (e \mu \nu / \mu \mu \nu \\ \text{GGM } \tilde{g} \tilde{g}, \tilde{X}_1^{1} \rightarrow Z \tilde{G} \end{array}$	Disapp. trk dE/dx trk 0 trk dE/dx trk μ) $1-2\mu$ 2γ $displ. ee/e\mu/\mu$ displ. vtx + jet	1 jet - 1-5 jets - - - μ - is -	Yes Yes - - Yes - Yes	36.1 18.4 27.9 3.2 3.2 19.1 20.3 20.3 20.3	$\begin{array}{c c} \dot{x}_{1}^{\pm} & 430 \text{ GeV} \\ \ddot{x}_{1}^{\pm} & 495 \text{ GeV} \\ \ddot{g} & 850 \text{ GeV} \\ \ddot{g} \\ \dot{g} $	1.58 TeV 1.57 TeV	$\begin{split} &m(\tilde{\mathcal{E}}_1^+)-m(\tilde{\mathcal{E}}_1^0)-160\;MeV,\;\tau(\tilde{\mathcal{E}}_1^+)=0.2\;ns\\ &m(\tilde{\mathcal{E}}_1^+)-m(\tilde{\mathcal{E}}_1^0)-160\;MeV,\;\tau(\tilde{\mathcal{E}}_1^+)<15\;ns\\ &m(\tilde{\mathcal{E}}_1^0)=100\;GeV,\;10\;\mus\!<\!\tau(\tilde{\varrho})\!<\!1000\;s\\ &m(\tilde{\mathcal{E}}_1^0)\!=\!100\;GeV,\;\tau\!>\!10\;ns\\ &10\!\cdot\!tan_{\mathcal{B}}\!<\!50\\ &10\!\cdot\!\tau(\tilde{\mathcal{E}}_1^0)\!<\!3\;ns,\;SPS8\;model\\ &1\!<\!\tau(\tilde{\mathcal{E}}_1^0)\!\!<\!\!740\;mn,\;m(\tilde{\varrho})\!=\!\!1.3\;TeV\\ &6\!<\!c\tau(\tilde{\mathcal{E}}_1^0)\!<\!480\;mn,\;m(\tilde{\varrho})\!=\!\!1.1\;TeV \end{split}$	ATLAS-CONF-2017-017 1506.05332 1310.6584 1606.05129 1604.04520 1411.6795 1409.5542 1504.05162
RPV	$ \begin{array}{l} LFV \ p + X, \widetilde{v}_{\tau} \rightarrow e \mu / e \tau / \mu \tau \\ Bilinear \ RPV \ CMSSM \\ \widetilde{X}_1^{\dagger} \widetilde{X}_1^{-}, \widetilde{X}_1^{+} \rightarrow W \widetilde{X}_1^{0}, \widetilde{X}_1^{0} \rightarrow e e v, e \mu v, \mu \mu \\ \widetilde{X}_1^{\dagger} \widetilde{X}_1^{-}, \widetilde{X}_1^{+} \rightarrow W \widetilde{X}_1^{0}, \widetilde{X}_1^{0} \rightarrow \tau \tau v_e, e \tau v_\tau \\ \widetilde{g} \widetilde{g}, \widetilde{g} \rightarrow q q \\ \widetilde{g} \widetilde{g}, \widetilde{g} \rightarrow q q \\ \widetilde{g} \widetilde{g}, \widetilde{g} \rightarrow q \widetilde{\chi}_1^{0}, \widetilde{\chi}_1^{0} \rightarrow q q q \\ \widetilde{g} \widetilde{g}, \widetilde{g} \rightarrow \widetilde{t} \widetilde{\chi}_1, \widetilde{\chi}_1^{-} \rightarrow q q q \\ \widetilde{g} \widetilde{g}, \widetilde{g} \rightarrow \widetilde{t} \widetilde{t}, \widetilde{t}_1 \rightarrow b s \\ \widetilde{t}_1 \widetilde{t}_1, \widetilde{t}_1 \rightarrow b s \\ \widetilde{t}_1 \widetilde{t}_1, \widetilde{t}_1 \rightarrow b \ell \end{array} $	$\begin{array}{c} e\mu, e\tau, \mu\tau \\ 2 e, \mu \ (SS) \\ y & 4 e, \mu \\ 3 e, \mu + \tau \\ 0 & 4 \\ 1 e, \mu & 8 \\ 1 e, \mu & 8 \\ 0 \\ 2 e, \mu \end{array}$	- 0-3 <i>b</i> - - 5 large- <i>R</i> jel 5 large- <i>R</i> jel -10 jets/0-4 -10 jets/0-4 2 jets + 2 <i>b</i> 2 <i>b</i>	Yes Yes Yes ts - ts - b - b -	3.2 20.3 13.3 20.3 14.8 14.8 36.1 36.1 15.4 36.1	\$\vec{v}_r\$ 1 \$\vec{u}_i\$ 1.14 TeV \$\vec{u}_1\$ 1.14 TeV \$\vec{u}_1\$ 450 GeV \$\vec{u}_2\$ 1.08 TeV \$\vec{u}_2\$ 1.0	1.9 TeV 1.45 TeV V 1.55 TeV 1.55 TeV 1.65 TeV 1.65 TeV	$ \begin{split} & \lambda_{311}^{l} = 0.11, \lambda_{132/133/233} = 0.07 \\ & m(\tilde{q}) = m(\tilde{g}), c\tau_{LSP} < 1 \text{ mm} \\ & m(\tilde{x}_{1}^{0}) > 400 \text{GeV}, \lambda_{122} \neq 0 (k = 1, 2) \\ & m(\tilde{x}_{1}^{0}) > 0.2 \times m(\tilde{x}_{1}^{\dagger}), \lambda_{133} \neq 0 \\ & \text{B}(r) = \text{B}R(b) = \text{B}R(c) = 0\% \\ & m(\tilde{x}_{1}^{0}) = 800 \text{ GeV} \\ & m(\tilde{x}_{1}^{0}) = 1 \text{ TeV}, \lambda_{112} \neq 0 \\ & m(\tilde{r}_{1}) = 1 \text{ TeV}, \lambda_{323} \neq 0 \\ & \text{B}R(\tilde{r}_{1} \rightarrow be/\mu) > 20\% \end{split} $	1607.08079 1404.2500 ATLAS-CONF-2016-075 1405.5086 ATLAS-CONF-2016-057 ATLAS-CONF-2016-057 ATLAS-CONF-2017-013 ATLAS-CONF-2017-013 ATLAS-CONF-2017-013 ATLAS-CONF-2016-084 ATLAS-CONF-2017-036
Other	Scalar charm, $\tilde{c} \rightarrow c \tilde{\chi}_1^0$	0	2 c	Yes	20.3	č 510 GeV		$m(\tilde{\chi}_1^0)$ <200 GeV	1501.01325

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*Only a selection of the available mass limits on new states or phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

D Cl.

10⁻¹

Mass scale [TeV]

The future





ATLAS Phase-I upgrades

LHC luminosity (collision rate) has exceeded LHC design already by 40%

Could exceed design luminosity by a factor ~2.5 in "Run-3"

→ Phase-1 upgrades give us better first-level trigger performance (better selectivity in hardware within ~3 μ s), and also to provide better tracking close to the interaction point

Main ATLAS Phase-I upgrades:

- New inner pixel layer installed already in 2014
 - Having impact e.g. for $H \rightarrow bb$ analysis
- New track & calorimeter trigger electronics
- New "small muon wheel" (9.3m diameter)







Protons physics Commissioning

Ions

HL-LHC 14 TeV-





HL-LHC accelerator upgrade was approved by CERN Council in June 2016 (cost: 930M CHF) HL-LHC: "levelled" luminosity 5-7 times the original design, until ~2035

Accumulate 10x more data than in Runs 1-3 combined - era of high precision

Must upgrade detectors!



HL-LHC: precision Higgs physics

All measurements in Higgs sector will benefit strongly from more statistics, right through HL-LHC programme

- Start measuring rarer decays (SM, BSM?)
- Higher precision measurements
 - Probe Higgs couplings at ~percent level to look for sign of new physics in loops
- Search for anomalous production etc etc





HL-LHC: di-Higgs production



The Higgs potential V is fundamental to the "Brout-Englert-Higgs mechanism" → non-zero scalar field *in vacuo*

Essential to explore shape of the potential - beyond quadratic term $\rightarrow m_{_{H}}$ \rightarrow also to throw light on electroweak phase-transition in early universe



HL-LHC: di-Higgs production



H pair production provides this sensitivity, in principle Range of studies going on to assess λ sensitivity

- Many channels
- High backgrounds will need to be measured from the data





Indications we can measure HH rate at ~30-50% level with the full HL-LHC data sample - maybe better \rightarrow sensitivity to anomalous large λ

Studies continue...

HL-LHC: extended search reach

With ten times the luminosity, mass reach for new particles increases by roughly 30% for many processes

For new processes with lower cross-sections, the sensitivity gain can be much higher

• e.g. weakly-coupled dark matter models

Example shown: Electroweak production of SUSY partners of electroweak bosons







Current ATLAS







D Charlton / Birmingham - October 2017 - Beyond the Higgs Boson Discovery

Phase-II tracking detector: ITk



Inner Tracker Strip

All silicon-sensor tracker:

- inner layers pixel sensors
- outer layers strip sensors

Sensor and systems R&D ongoing for some years



Closing Run: 191426 Event: 86694 2011-10-22 1



The LHC is delivering well-beyond-design luminosity, at 13 TeV pp collision energy in Run-2

- Rich and diverse physics programme (also LHCb, ions)
- Many searches for new particles and interactions in progress
- Irrespective of new discoveries, there is a broad precision physics programme at ATLAS and CMS

The Higgs boson discovery was just the start of the programme of study of the scalar sector at the LHC

- We know now that it is a Higgs scalar, with broadly SM-like properties
- Huge scope for new physics in the scalar sector we have just scratched the surface so far
- Only the LHC will address these questions for the next two decades.

The LHC is, and will remain to the mid-2030's, the world's energy frontier particle collider