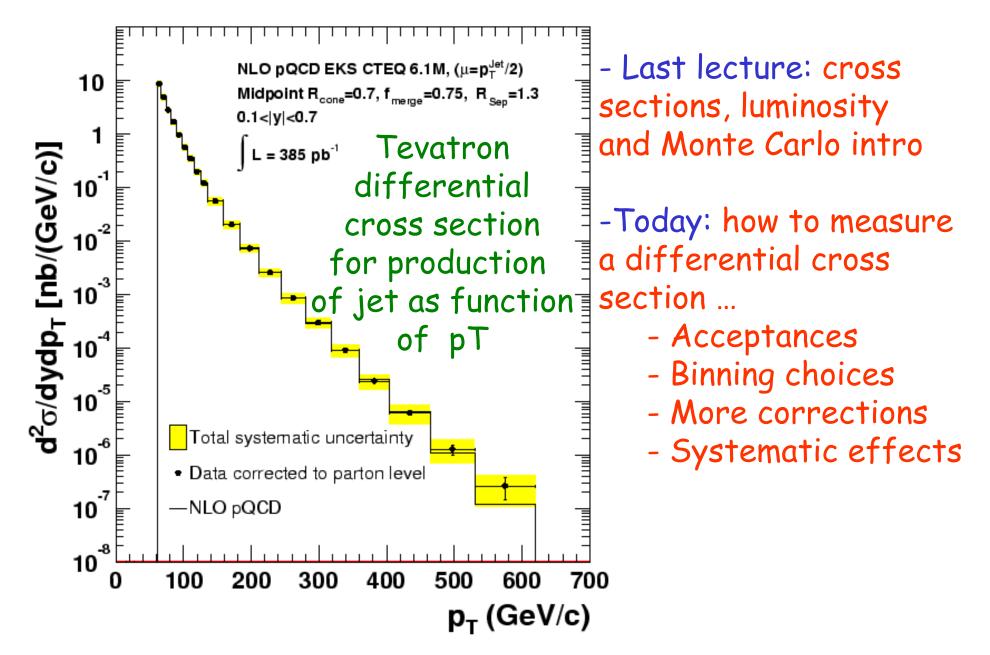
Binned Cross Section Measurements



Reminder: Cross Sections

A crucial link between experimental and theoretical worlds!

$$\sigma = \frac{N_{obs} - N_{b/g}}{L \cdot A_{cc}}$$
 ... where efficiency ϵ absorbed in A_{cc}

• A cross section should be uniquely defined in terms of configurations of <u>observable</u> particles... i.e. in terms of observed hadrons, not Partons or `Leading Order'!...

theories develop and improve; data are fixed by nature!

• Definitons in terms of **Lorentz invariants** preferred!

 $N_{b/g}$ and A_{cc} can be obtained from Monte Carlo models, but should always check that our analysis is <u>ROBUST</u> ... i.e. relatively insensitive to the details of the Monte Carlo!!!

Reminder: Monte Carlo Simulations

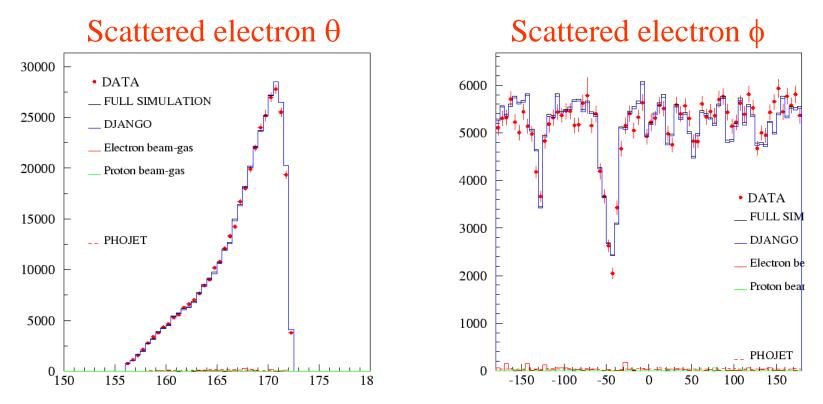
- 3 steps: 1) Event <u>generator</u> models underlying physics
 ... outputs a list of `stable' particles
 - 2) <u>Simulation</u> of detector response to particles
 - 3) **<u>Reconstruction</u>** identical to that of data
- If all steps are done well, MC output should closely model the data, but crucially, we have full info about parton and hadron levels as well as reconstructed level!

... Can answer questions about our measurement. e.g:

- What was my acceptance and efficiency?

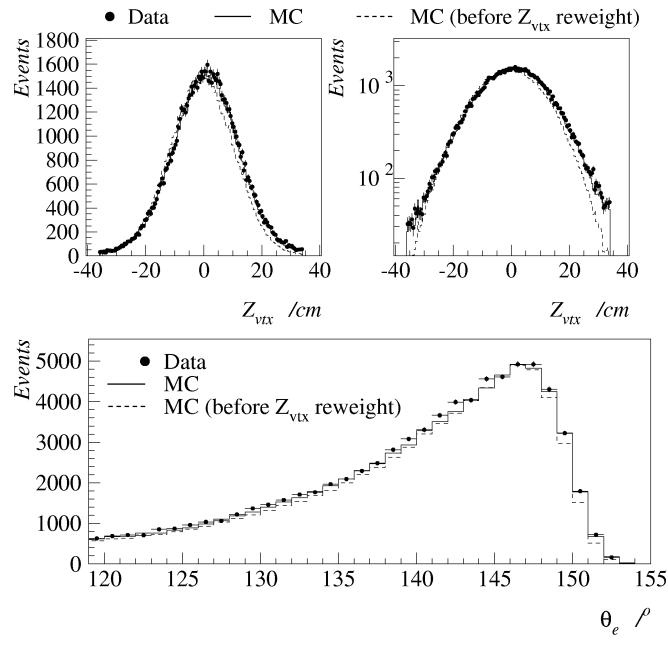
- What fraction of my `electrons' are really photons? Answers are (slightly) affected if underlying generator is poor model of physics. Test with `control distributions' e.g. Control Distributions of reconstructed data (H1)

... compare data before any corrections with rec level MC ...



- Good description by DJANGO DIS Monte Carlo simulation ③
- Big holes in ϕ due to dead bits of detector but described by simulation
- Background Monte Carlo (PHOJET) gives tiny contribution
- Reassuring results happy with model use to correct for inefficiencies such as the holes!

Bad Control Distributions and their Influence



In this H1 example, distribution of simulated z position of interaction vertex is slightly wrong due to a mistake in steering parameters

Fixing vertex distribution results in much better description of an important phsysics distribution – electron scattering angle.

What is a <u>Differential</u> Cross Section

Often we are interested in the dependence of a cross section on a variable (e.g. transverse energy of a jet, pseudorapidity, momentum transferred ...)

Work in terms of `differential cross sections e.g. d σ / dE_{t}

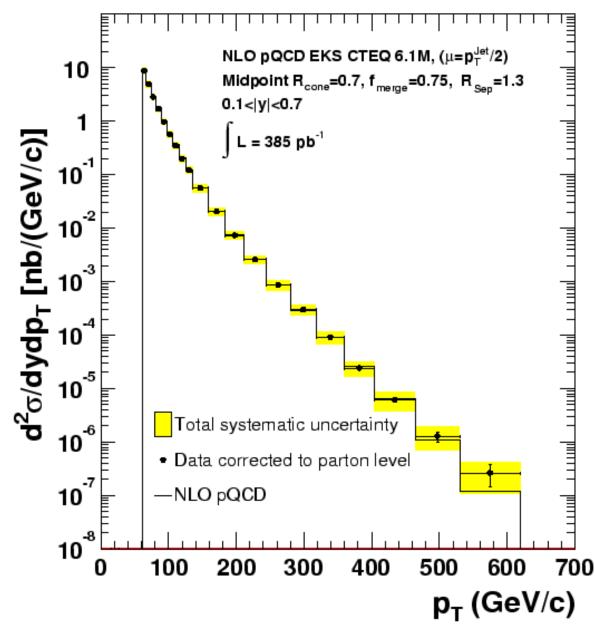
Definition:
$$\frac{d\sigma}{dx} = \lim_{\Delta x \to 0} \frac{\Delta \sigma}{\Delta x}$$
 such that $\int_{-\infty}^{\infty} \frac{d\sigma}{dx} dx = \sigma_{tot}$

... measure a differential cross section at a single fixed point in x by creating a bin of width Δx about that point ... modification to cross section formula ...

$$\frac{d\sigma}{dx} = \frac{N_{obs} - N_{b/g}}{L \cdot A_{cc} \cdot \Delta x}$$

... where Δx is the width of our bin in x (still have to worry about exactly which x value we measured at)

An Example Differential Cross Section



In this example, a jet cross section is measured differentially in the jet transverse momentum p_T and the jet rapidity y

... note the horizontal error bars indicating that the data point could correspond to any point in the bin ... (pessamistic!!!)

`Generalised' or `Smeared' Acceptance

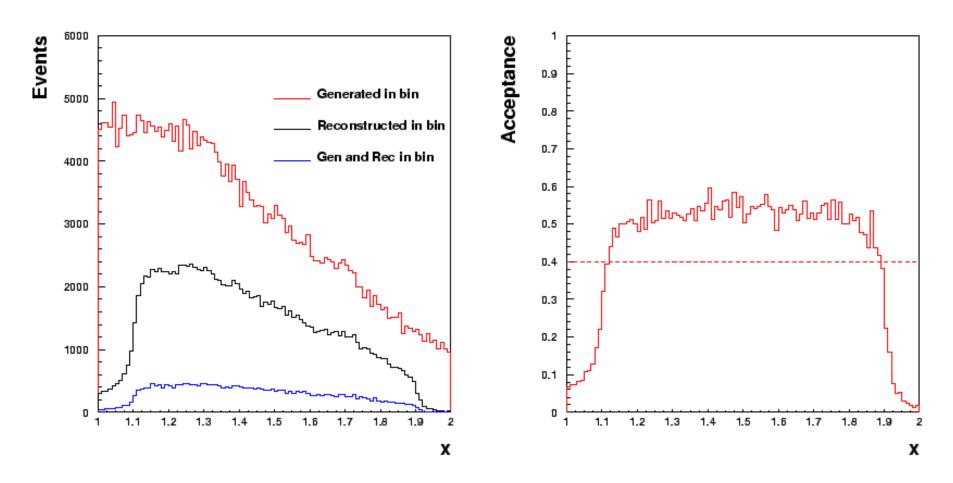
Introduce smeared acceptance A_{cc} Recall: $\sigma = \frac{N_{obs} - N_{b/g}}{L \cdot \epsilon \cdot A_{cc}}$ (calculated from Monte Carlo), which corrects for 'everything' in binned cross section measurements:

- 1. Finite selection efficiency / acceptance: some events end up in no bin at all.
- 2. Finite resolution: events end up in wrong bin (migrations).

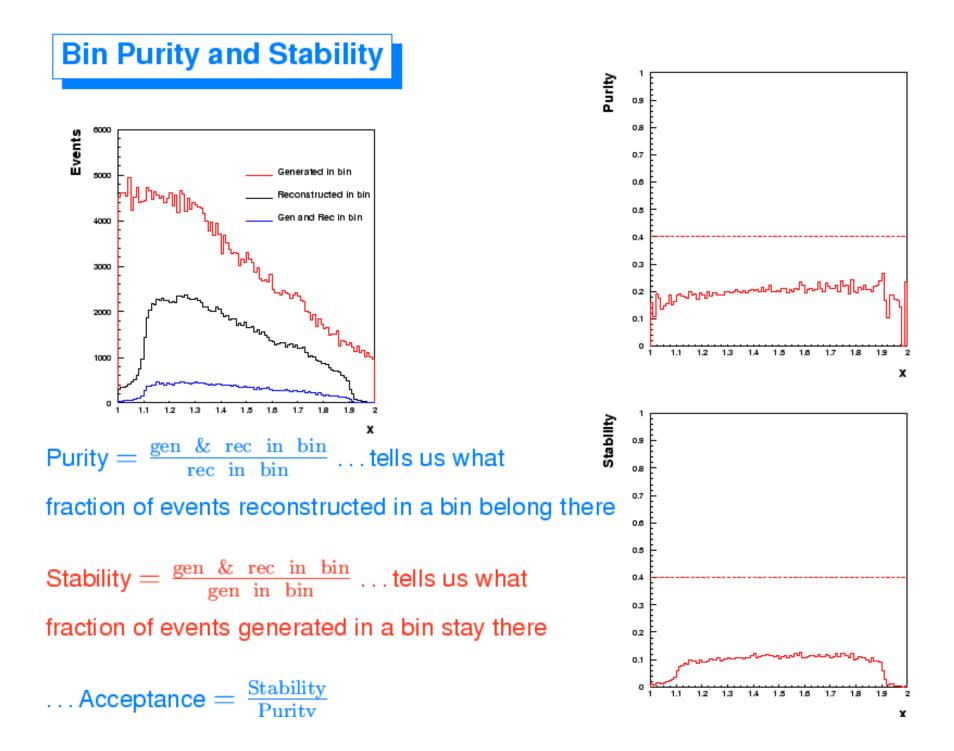
Calculating A_{cc} with an MC can be very easy!... $A_{cc} = N_{rec} / N_{aen}$

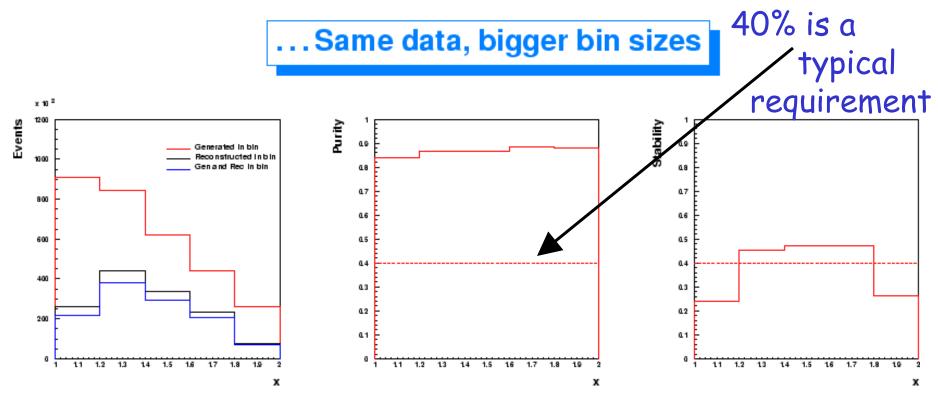
This simple treatment is OK if resolutions are well behaved (Gaussian) and bin choices are sensible (migrations not too large) ... otherwise more complex 'unfolding' needed.

A Real Acceptance Example



Here acceptance $\sim 50\%$ for $1.1 \leq x \leq 1.9 \dots \rightarrow$ can make measurement ... but is acceptance the full story? ... here very few events are reconstructed and generated in same bin ... we need many fewer bins!

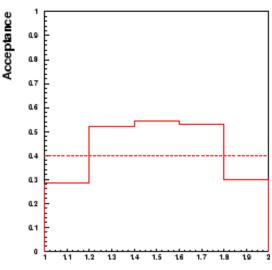




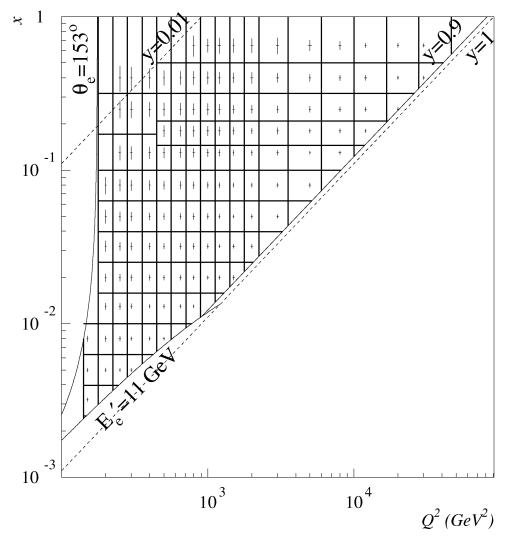
With bigger bin sizes, purity and stability improve whereas acceptance is unchanged.

Measurement region is determined by acceptance ... must be high and stable across bin

Binning choices usually defined by purity e.g. require purity > 0.68 ... i.e. bin width > resolution Exception: if data statistics are limited



e.g. Binning Choice from H1 (double differential)



Example from a 2 dimensional measurement (x and Q² in DIS)

Crosses indicate resolution calculated in each bin, always smaller than bin size

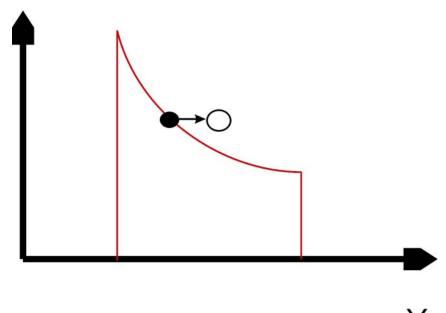
Optimal measurement required variable bin sizes due to resolution and available statistics (decreasing as Q² increases)

All acceptances, purities and stabilities (as calculated from Monte Carlo models) are well above 40% ...

More corrections with MC: Finite Bin Size Effects

The cross section we measure is for the `centre of gravity' of the bin.
For a cross section which falls fast, this can be very different from the cross section at the bin centre.

Plotting at the bin centre without correction is a <u>very</u> common mistake
One solution is to show horizontal error bars (as in Tevatron example)



Х

Better to plot at the bin centre, x_c , and make a **`bin centre correction'** back to the Monte Carlo again to calculate as ...

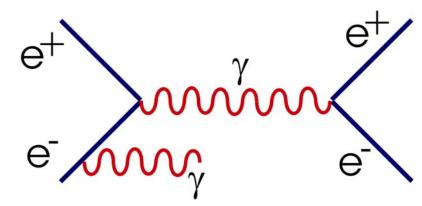
BCC = Cross section at bin centre

Cross section averaged over bin

Using MC to deal with Radiative Effects

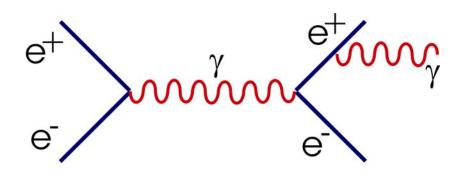
- For many measurements we want to correct for higher order QED effects such as initial state photon radiation.
- Whether we do this or not is a question of definition, but theorists probably don't want to have the radiative effects included.
- They want to work with a **BORN Level Cross Section**

1) Initial state radiation (ISR) effectively lowers beam energy of some events



Photon emitted usually at very small angle to beam electron – sometimes detected e.g. in small angle calorimeter but usually not

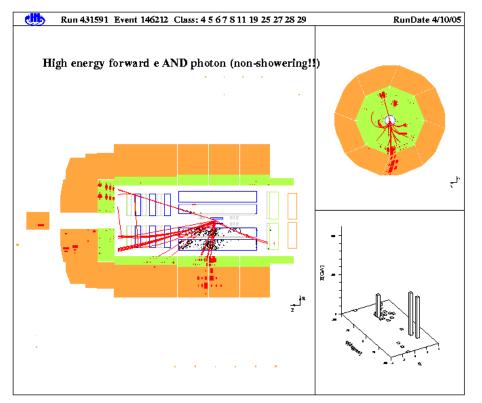
Radiative Effects II: Final State Radiation



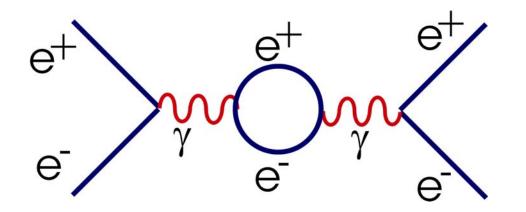
• As for ISR, final state photon radiation is usually nearly collinear with radiating electron – in fact you will usually reconstruct it in the same calorimeter cluster – in which case no correction needed!

• Still need small corrections for wide angle emissions

ep example



Radiative Effects III



Third effects is `virtual loop corrections' – effectively just a change in photon coupling

Monte Carlo models to the rescue again, as long as it has an interface to a reliable QED model! Correct for all 3 effects with a factor

Predicted cross section with ISR, FSR and loops switched off in MC

Predicted cross section with ISR, FSR and loops simulated in MC

Radiative corrections can be large e.g. if measurement involves photons or is very sensitive to the beam energy

What about the Uncertainties?

• We now have all ingredients needed to obtain the basic result for a (binned) differential cross section ©

• However, this is not really a `*measurement*' until we also assess the errors ... which is usually most of the work \otimes

• Future lectures will deal with statistical errors (always calculable in principle, but often done incorrectly) and systematic errors (generally a poorly understood subject, even by professionals, but still almost[?] always quantifiable in principle

... first, some basic error theory, which applies to both ...

Reminder of Mean, Standard Deviation ...

Mean / Average of a distribution

 $\mu(x)=\langle x
angle=~ar{x}=\sum_{m{x}}m{p}_{m{x}}\cdotm{x}$... or $ar{x}=\int P(x)\cdot xdx$ if continuous distribution

Mean of a derived quantity

$$\overline{f(x)} = \sum_x p_x \cdot f(x)$$
 e.g. If $f(x) = 2x$, $\overline{f(x)} = \sum_x p_x \cdot 2x = 2\overline{x}$

Variance V

Average squared deviation from mean used for spread (average deviation from mean is zero)

$$V(x) = \overline{(x-\overline{x})^2} = \sum_x p_x (x-\overline{x})^2 = \overline{x^2} - (\overline{x})^2$$

Standard Deviation

Often used in preference to Variance because it has the same units as $m{x}$

$$\sigma(x) = \sqrt{V(x)} \qquad = \sqrt{rac{1}{N}\Sigma(x-ar{x})^2}$$

• For most purposes, assume that measurements follow a Gaussian distribution about the true value of a variable. • By the `error' on a measurement, we mean 1σ (Gaussian)